

Electrical impedance of the hydrophobic oil-water interface: contactless measurement

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Abstract: We have performed systematic contactless measurements aiming at the determination of the electrical impedance of an oil-water interface. The oil-water interface was formed within an insulin syringe tube due to the oil hydrophobicity. The measurement AC signal (0.1-300 kHz at 0.1 V) was capacitively coupled to the oil-water system through the syringe wall. The interface resistivity was estimated at $\sim 2 \times 10^9 \Omega \cdot \text{m}$ (assuming the interface as a slip layer with width ~ 10 nm) and was associated with hydroxide (OH^-) ions, agreeing with existing reports. The resistivity was found to be independent on the signal frequency, which may indicate the nonpolarized nature of the hydrophobic interface. The present impedance spectroscopy is important for hydrophobic systems and may give an insight into the future nondestructive (contactless) investigation of the electrical properties of the lipid layer in the living cell membrane.

Keywords: hydrophobicity, Oil-water interface, Impedance spectroscopy, contactless measurement.

1. Introduction

Hydrophobic materials (such as oil) are non-polar and, thus, have a weak affinity to water. Applications of hydrophobic materials are diverse [1] and include the renowned superhydrophobic self-cleaning surfaces [2] and anti-icing surfaces [3]. The hydrophobicity effect has a direct impact on many chemicals (e.g. [4]), biological (e.g. [5]), and environmental (e.g. [6]) processes. Recently, the incorporation of materials that are hydrophobic (such as the reduced graphene oxide) and hydrophilic (a material with a strong affinity to water such as 2D graphene oxide) in a gas sensor has enhanced its sensing capability to overcome a challenging issue of distinguishing between tetrahydrofuran and dichloromethane [7].

These applications have stimulated fundamental research aiming at the investigation of the electrical properties of the hydrophobic interface. For instance, Tian *et al* [8] managed to experimentally find that the hydrophobic interface is negatively charged due to the high adsorption of the hydroxide (OH^-) ions at the interface. This finding has been supported by molecular dynamics simulations [9]. Hydrophobic materials have also been utilized for controlling the metallic corrosion rate (e.g. [10] and [11]). The results of impedance spectroscopy measurements on Cu with hydrophobic coating have shown a significant reduction in the corrosion rate in seawater [12].

To contribute to the ongoing fundamental research on the electrical characterization of the hydrophobic interface, this paper reports on the systematic contactless measurements of the impedance of the oil-water hydrophobic interface. We have used commercial olive oil and distilled water within an insulin syringe. The oil-water interface was formed due to the oil hydrophobicity. The measurement AC signal was capacitively coupled to the oil-water system through the syringe plastic wall. We have previously used a similar contactless measurement technique to accurately determine the dependence of the conductivity on the size of gold nanoparticles [13]. In the present work, we used the Solarton Analytical 1287A and 1255B (see Fig. 1) Frequency Response Analyzer [14] which has good measurement accuracy and precision. The reactance of the plastic syringe wall $< 100 \text{ k}\Omega$ at the measurement frequencies of 100 kHz $< f < 300$ kHz, which allows sufficient portion of the measurement signal to penetrate the plastic wall into the oil-water interface. Hence, we have managed to estimate the resistivity of the interface at $\sim 2 \times 10^9 \Omega \cdot \text{m}$ and found it to be consistent with ionic resistivity associated with OH^- in water, a result that conforms with existing investigations (e.g. [9]) that attribute the negative charge of the interface to the OH^- ions.

Besides being nondestructive and contamination-free, contactless impedance measurement has become an important method for the detection of inorganic or small organic ions in capillary electrophoresis [15]. Thus, we anticipate that our present work would open an interesting research avenue towards thenondestructive (contactless) investigation of hydrophobicity associated with a lipid layer (e.g. [16] and [17]) in the living cell membrane.

2. Measurement procedure and impedance model

Figure 1 shows a schematic of the measurement set-up consisting of the measuring instrument (Solarton Analytical 1287A and 1255B Frequency Response Analyzer [14]) from which the AC signal (100 mV, 0.1 – 300 kHz) is applied to the pair of Aluminum electrodes (Al electrodes) that are tightly wrapped around the syringe tube (inner diameter ~ 2 mm). Thus, the electrodes are capacitively coupled to the sample (oil-water) within the tube through the syringe plastic wall. We use commercial olive oil and distilled water. The effective length of the channel within the tube containing the sample is defined by the electrode-electrode separation which was $l \sim 1.5$ cm. The hydrophobic interface prevents oil-water mixing.

The top part of Fig. 1 is the impedance model for the sample and the surrounding (external) electrical environment. The impedances $Z_w, Z_o, Z_i, Z_p,$ and Z_x are for the water, oil, oil-water hydrophobic interface, tube wall, and external circuit (including parasitic ones). From the model, the effective measured impedances for a channel with water only, oil only, and oil-water samples are respectively given by equations (1) – (3) below:

$$Z_{mw}^{-1} = Z_x^{-1} + (Z_p + Z_w)^{-1} \dots\dots\dots (1)$$

$$Z_{mo}^{-1} = Z_x^{-1} + (Z_p + Z_o)^{-1} \dots\dots\dots (2)$$

$$Z_{mow}^{-1} = Z_x^{-1} + \left(Z_p + \frac{Z_o}{2} + Z_i + \frac{Z_w}{2} \right)^{-1} \dots\dots\dots (3)$$

In equation (3), Z_o and Z_w are divided by 2 because the two halves of the channel were filled with oil and water. From eqns. (1) – (3) we deduce:

$$Z_i \approx (Z_{mow}^{-1} - Z_{ma}^{-1})^{-1} - \frac{(Z_{mw}^{-1} - Z_{ma}^{-1})^{-1}}{2} - \frac{(Z_{mo}^{-1} - Z_{ma}^{-1})^{-1}}{2} \dots\dots\dots (4)$$

Where $Z_{ma} \approx Z_x$ is the measured impedance with empty tube (i.e. air-filled tube). This approximation holds because the empty tube has a significantly large capacitive reactance, hence, the measurement effectively gave rise to Z_x .

It is worth mentioning that the measured impedances in equation (4) were provided in terms of their measured real and imaginary parts. Therefore, a simple MATLAB® or Microsoft® Excel code was used to determine Z_i from the complex-number analysis of equation (4).

3. Results and discussions

Figure 2 (a) and 2(b) show the measured impedances $Z_{ma}, Z_{mw},$ and Z_{mow} versus the frequency, f . To avoid crowding the space, we have not shown Z_{mo} because we observed that $Z_{mo} \sim Z_{ma}$, indicating high oil impedance due to lack of polarity. Clearly, the data is complicated due to the incorporation of the unwanted (parasitic) external impedance, Z_x . However, as shown in section 2, we use the impedance model analysis to eliminate Z_x and extract the impedance of interest from the measurement data. For example, Z_i (oil-water interface impedance) was extracted in equation (4) and its real and imaginary parts (**Re. Z_i and $-\text{Im}Z_i$**) are illustrated in Fig. 2(c) as a function of the frequency. Clearly, the real part (open blue triangles) is independent on the frequency (f) as can be seen from the red line which is shown as a guide for the eye. This suggests that the oil-water hydrophobic interface is dominantly resistive specifically because the imaginary part (open dark squares) did not follow the $1/f$ capacitive behaviour as, for instance, seen for the imaginary parts in Fig. 2(a) and 2(b). The $\text{Im}Z_i$ may thus be associated with the residual of the calculation leading to equation (4). Such residual impedance may arise from the approximation $Z_{ma} \sim Z_x$ and some hidden stray impedances. We observed similar residual imaginary part in a measurement of known resistances conducted to verify the accuracy of our model and the reliability of the measurement setup. This measurement will be discussed in conjunction with Fig. 3. We will show that the known resistances were measured with good accuracy regardless of the presence of this residual impedance.

Nevertheless, even if the observed $\text{Im}.Z_i$ was physical, then Z_i may roughly resemble an ideal supercapacitor [18] that may be modelled [19] as a resistance, R , in series with a constant-phase-element, CPE [20] i.e. an imperfect capacitor. In this case, the impedance is characterized by a frequency-independent real part representing R , in an agreement with our data (Fig. 2(c)). However, we have not investigated the possibility of this series R-CPE model because the expected $1/\omega$ behaviour of $\text{Im}.Z_i$ was obscured by the noisy $f < 100$ kHz region as seen in Fig. 2(c).

We stress that we have not observed signature of carrier diffusion across the oil-water interface that can be modelled using, for instance, Warburg diffusion elements [21].

In addition to the above-mentioned technical justifications, our assumption that the interface is dominantly resistive may be justified and supported by the molecular dynamic results by Kudinet al [9] which confirmed that the hydrophobic interface (between water and graphene) has a net negative charge, i.e. effectively nonpolar, as implied by the experiment (e.g. [8]). The absence of charge polarity at the interface may crudely be regarded as a signature of a resistive interface that would otherwise exhibit a pronounced capacitive and, hence, frequency dependent contribution to Z_i . To arrive at a more decisive conclusion about the resistive nature of the oil-water interface, we need to construct and analyze a sophisticated model of a system of a charged interface between polar electrolyte (representing the water) and insulator (representing the oil). A similar model was the subject of recent research [22] on double-layer (DL) capacitance [23] which showed that the capacitance was independent of the system size. We assume that the change in system size corresponds to change in charge density, Q , at the interface with applied potential, ξ , from the AC signal. Thus, the interface capacitance $\sim dQ/d\xi \sim \text{constant}$, resembling the equilibrium differential capacitance regime for the DL illustrated in Fig. 1 of Ref. # [24] with constant overall resistance. Our data in Fig. 2(c) is characterized by the constant $\text{Re}.Z_i$ similar to the constant overall resistance in Ref. # [24]. Thus, we associate $\text{Re}.Z_i$ with the interface resistance as follows.

We assigned the resistance of the hydrophobic interface to the real part of Z_i , $\text{Re}.Z_i$. From Fig. 2(c), $\text{Re}.Z_i \sim 6 \text{ M}\Omega$. The resistivity of the interface was estimated as $\rho_i \sim (\text{Re}.Z_i) \times A / t \sim 2 \times 10^9 \Omega \cdot \text{m}$, where $A \sim 0.03 \text{ cm}^2$ is the cross-sectional area of the interface, t is the effective thickness of the interface assumed to be similar to the slip length $\sim 10 \text{ nm}$ as implied by Zhenbin et al [25] and Maali et al [26]. Kudinet al [9] attributed the negative interface charge to the affinity of the negative hydroxide ions (OH^-) to the hydrophobic interface at low pH. We used commercial distilled water that has $\text{pH} \sim 5$ to 7 . We assume that the interface has roughly the same OH^- concentration as that for the water ($\sim 10^{-14}/10^{-\text{pH}}$) which ranges from 10^{-9} to 10^{-7} molar. Consequently, the corresponding OH^- conductivities are between $\sim 10^{-8} \Omega^{-1} \text{ m}^{-1}$ and $10^{-6} \Omega^{-1} \text{ m}^{-1}$ which were roughly estimated using the OH^- molar conductivity of $19.8 \times 10^{-3} \Omega^{-1} \text{ m}^2 \text{ mol}^{-1}$ [27]. According to our measurement, the hydrophobic interface conductivity $= 1/\rho_i \sim 10^{-9} \Omega^{-1} \text{ m}^{-1}$ which is consistent with the estimated range of conductivities for OH^- in water.

To check the reliability of our measurement and impedance models, we analyzed the model for Z_{mw} measurement. The objective is to estimate the capacitance of the syringe plastic wall, C_p , then compare it with the theoretical value $\sim 4 \text{ pF}$ for the cylindrical plastic wall (0.2 mm thick with dielectric constant ~ 3) of the syringe beneath the two electrodes (each $\sim 2 \text{ mm}$ wide). From equation (3) and $Z_{ma} \sim Z_x$ we get $(Z_{mw}^{-1} - Z_{ma}^{-1})^{-1} = Z_p + Z_w = Z_{pw}$ which represents a series combination of Z_p (considered to be purely capacitive with capacitance C_p) and Z_w (considered as a parallel resistor (R_w)-capacitor (C_w) configuration (inset Fig. 2(d)). The main Fig. 2(d) shows that the real part, $\text{Re}.Z_{pw}$, vs. the imaginary part, $\text{Im}.Z_{pw}$, are nicely fit (blue and red lines, with impedance Z_{fit}) to the circuit model of Z_{pw} shown in the inset of Fig. 2(d). From the fitting, $C_p \sim 5 \text{ pF}$, $R_w \sim 0.2 \text{ M}\Omega$, and $C_w \sim 10 \text{ pF}$. Clearly, C_p from the model fitting agrees very well with the theoretical value $\sim 4 \text{ pF}$. The impedance analysis software used in Fig. 2(d) was the Matlab[®] Z-fit code [28].

To furthermore check the reliability of our model analysis (section 2) and measurement procedure, we used our measurement set-up to measure known resistances ($R_k = 0.47, 47, \text{ or } 94 \text{ k}\Omega$, inset of Fig. 3) of resistors impeded between two impedances (Z_{left} and Z_{right} , inset of Fig. 3). This measurement aimed to emulate the extraction method of Z_i (hydrophobic interface impedance) discussed in section 2. Here, R_k , Z_{left} , and Z_{right} (inset of Fig. 3) played the roles of Z_i , Z_o and Z_w (inset of Fig. 1), respectively. The results were then analyzed following similar procedures as that presented in section 2, leading to the following equation (5) that resembles equation (4):

$$R_k \approx (Z_{mlr}^{-1} - Z_{ma}^{-1})^{-1} - (Z_{ml}^{-1} - Z_{ma}^{-1})^{-1} - (Z_{mr}^{-1} - Z_{ma}^{-1})^{-1} \dots \dots \dots (5)$$

Where $Z_{ma} \sim Z_x$, and the measured impedances Z_{mlr} , Z_{ml} , and Z_{mr} , respectively, correspond to measurements on Z_{left} - R_k - Z_{right} , Z_{left} only, and Z_{right} only. The main Fig. 3 shows the extracted R_k values (red, blue and dark squares) using equation (5) versus the frequency. Interestingly, the extracted R_k values were frequency independent and matched the known R_k values with good accuracy.

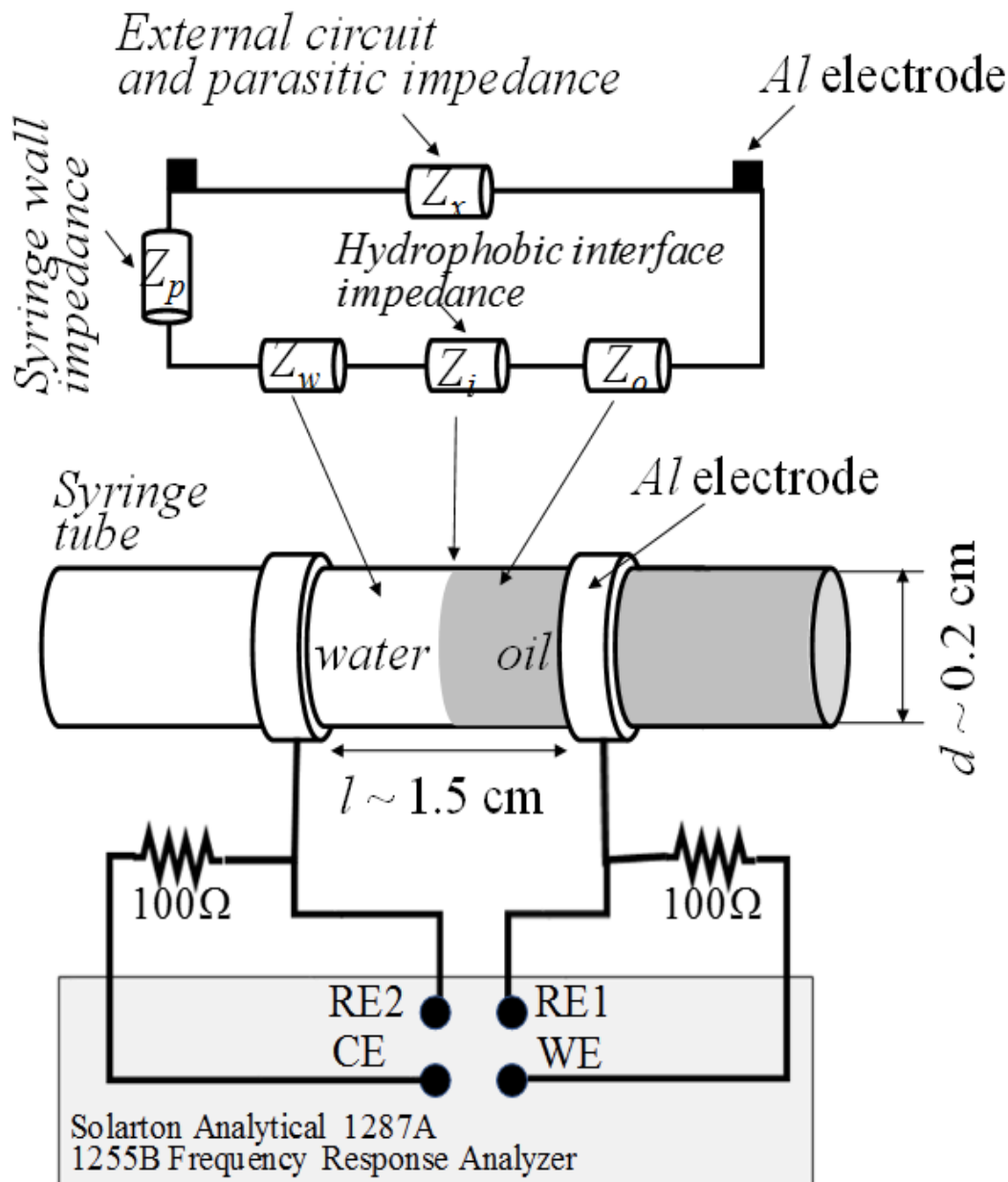


Figure 1: Schematic of the measurement set-up involving the impedance analyzer instrument (Solarton 1287A and 1255B). The sample (e.g. oil-water) is coupled capacitively to the Al electrodes through the syringe plastic wall. CE, WE, RE1 and RE2 are respectively the counter, working, reference 1, reference 2 electrodes of the instrument. The top circuit diagram depicts the impedance model of the measurement set-up where the impedances Z_w , Z_o , Z_i , Z_p , and Z_x are for the water, oil, oil-water hydrophobic interface, tube wall, and external circuit (including parasitic ones).

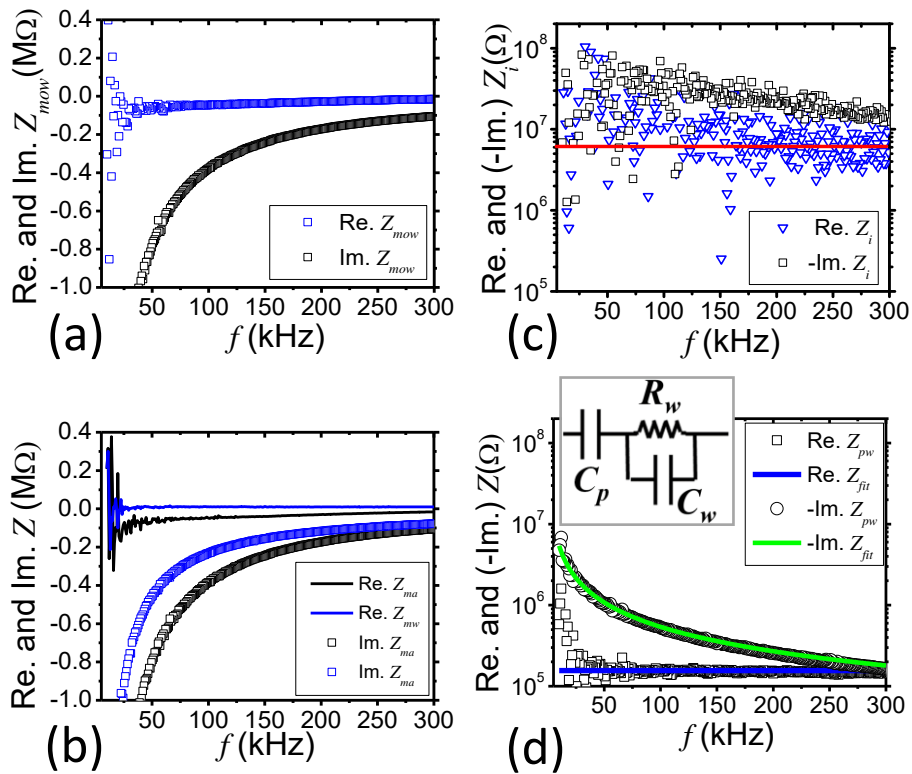


Figure 2: the measured impedances (Re. and Im. refer to real and imaginary impedance parts) versus the frequency, f , for (a) oil-water sample, Z_{mow} , and (b) empty tube, Z_{ma} , and water filled tube, Z_{mv} . (c) The hydrophobic interface impedance, Z_i , extracted using equation (4) (see text) (d) $-\text{Im. } Z_{pw}$ and $\text{Re. } Z_{pw}$ vs. f where Z_{pw} is the equivalent impedance of water impedance ($Z_w \equiv$ parallel R_w-C_w) in series with plastic tube impedance ($Z_p \equiv$ reactance with capacitance C_p). The green and blue lines are the fits (impedance $\equiv Z_{fit}$) of the data using the circuit model shown in the inset with $C_p = 5 \times 10^{-12}$ F, $C_w = 1 \times 10^{-11}$ F, and $R_w = 0.2$ M Ω .

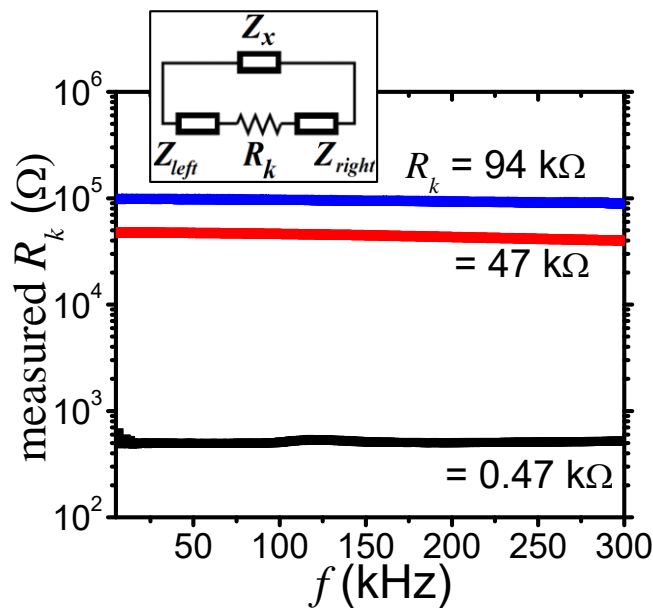


Figure 3: Measurement method of known resistances $R_k = 0.47, 47, \text{ or } 94 \text{ k}\Omega$ using the measurement setup of Fig. 1 where R_k were sandwiched in series between the impedances Z_{left} and Z_{right} (inset). The values of R_k (blue, red, black squares) extracted following similar model analysis as that presented in section 2 (see text).

4. Conclusion

Systematic contactless impedance measurements on oil-water interface have been performed at room temperature. The oil-water interface was formed within an insulin syringe tube due to the oil hydrophobicity. The measurement AC signal (0.1-300 kHz at 0.1 V) was capacitively coupled to the oil-water system through the syringe wall. The interface resistivity was estimated at $\sim 2 \times 10^9 \Omega \cdot \text{m}$ (assuming the interface as a slip layer with width $\sim 10 \text{ nm}$) from the measured equivalent impedances of the oil, water and empty tube and was found to be independent on the signal frequency, probably reflecting the nonpolarized nature of the hydrophobic interface. The present impedance spectroscopy is important for hydrophobic systems and may give an insight on future nondestructive (contactless) investigation of the electrical properties of the lipid layer in the living cell membrane.

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